

is probably possible as a result of the motion, thus re-establishing the relative equilibrium configuration.

A study which is based solely on intensities of diffraction maxima, gives only restricted information about the average equilibrium properties, and it is clear that much additional information is latent in the diffuse part of the scattering. It may be concluded at the present, however, that the molecule spends considerably more time at points close to the indicated equilibrium positions than elsewhere. A similar conclusion applies to the integrated time spent in these regions. Although the energy spectrum of the lattice is complex, the above conclusion indicates that the majority of the system is in quantum states lying between the extreme cases approximated by the harmonic oscillator and the free rotator (Eyring, Walter & Kimball, 1949) respectively, at temperatures a little below the melting point. This situation pertains also in a number of other cases which will be discussed elsewhere.

Consideration of the configurational entropies implied by the distinctly different statistical models discussed above; and comparison with the experimental values of the entropy changes at the low temperature phase transformations, may further illuminate the structural problems considered here.

The full-matrix least-squares refinement program is a somewhat revised version of the original form written by Gantzel, Sparks & Trueblood (I.U.C. World List No 384). The Fourier program was written by Gantzel & Hope (Department of Chemistry, University of California, Los Angeles 24, California).

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## Preparation, Structure and Properties of Triniobium Tetrasulphide, Nb<sub>3</sub>S<sub>4</sub>

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In the system niobium-sulphur a new phase, Nb<sub>3</sub>S<sub>4</sub>, has been prepared from the elements at 1000–1300°C. Fairly large needle-shaped crystals could be obtained by vapour transport with iodine. Nb<sub>3</sub>S<sub>4</sub> is hexagonal ( $a=9.5806$  Å;  $c=3.3747$  Å;  $c/a=0.35224$ ; space group  $P6_3/m$ ) and isotypic with Nb<sub>3</sub>Te<sub>4</sub> and Nb<sub>3</sub>Se<sub>4</sub>. The structure was refined from single-crystal diffractometer data obtained with Mo radiation. The structure is built up from NbS<sub>6</sub> octahedra which are linked together by shared faces and edges. Each metal atom is displaced by 0.306 Å from the octahedron centre towards one of the faces. In this way infinite zigzag metal chains are formed with Nb–Nb distances similar to those in Nb metal. Furthermore, the structure contains wide empty channels in the *c* direction. Nb<sub>3</sub>S<sub>4</sub> is a metallic conductor with temperature-independent paramagnetism.

### Introduction

The system Nb–S has been studied by Jellinek, Brauer & Müller (1960). The phases NbS<sub>3</sub>, 2*s*-NbS<sub>2</sub>, 3*s*-NbS<sub>2</sub>,

2*s*-Nb<sub>1+x</sub>S<sub>2</sub> and 3*s*-Nb<sub>1+x</sub>S<sub>2</sub> were characterized, but the phase relationships in the more metal-rich range could not be disentangled. An investigation of this range has, therefore, been undertaken. The new phases

found are  $\text{Nb}_3\text{S}_4$  and  $\text{NbS}_{1-x}$ , the latter existing in an orthorhombic high-temperature form and a hexagonal low-temperature form which have superstructures of the MnP and NiAs types respectively; the transition temperature is  $730^\circ\text{C}$  (Kadijk & Jellinek, 1968). The phase  $\text{Nb}_{21}\text{S}_8$  recently described by Franzen, Beineke & Conard (1968) was also observed in some of our samples. The present communication describes the preparation, physical properties and crystal structure of  $\text{Nb}_3\text{S}_4$ .

### Preparation and properties

Mixtures of niobium and sulphur in atomic proportions of about 1:1 were heated at  $1000\text{--}1300^\circ\text{C}$  in evacuated quartz tubes and cooled by quenching. Small needle-shaped crystals of a new phase were obtained, together with polycrystalline material. The crystals were hexagonal with  $a \approx 9.6 \text{ \AA}$ ,  $c \approx 3.4 \text{ \AA}$ ; the powder contained the same phase, together with several contaminants.

The shape of the hexagonal crystals and the unit-cell dimensions are similar to those reported for  $\text{Nb}_3\text{Te}_4$  and  $\text{Nb}_3\text{Se}_4$  (Selte & Kjekshus, 1964). Assuming the composition of the new phase to be  $\text{Nb}_3\text{S}_4$  and adopting the atomic parameters reported for  $\text{Nb}_3\text{Te}_4$ , qualitative agreement of observed and calculated powder intensities was obtained (Huisman, Kadijk & Jellinek, 1967). The composition  $\text{Nb}_3\text{S}_4$  was confirmed by a refinement of the crystal structure from single-crystal data.

Single crystals suitable for the X-ray intensity measurements were prepared by transport reactions using iodine as a carrier; various starting compositions and temperature gradients (about  $900^\circ \rightarrow 1000^\circ\text{C}$ ) were tried. In all cases two kinds of crystals grew in the high-temperature region of the quartz tube, namely hexagonal needles of  $\text{Nb}_3\text{S}_4$  and hard plates of orthorhombic  $\text{NbS}_{1-x}$ .

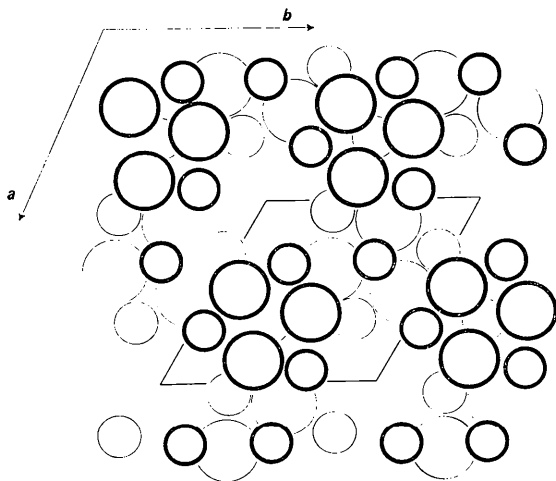


Fig. 1. The structure of  $\text{Nb}_3\text{S}_4$  projected along the  $c$  axis. Nb atoms are indicated by large circles, sulphur by small circles. Atoms at  $z = \frac{1}{4}$  are drawn with heavy lines, atoms at  $z = \frac{3}{4}$  with thin lines.

Some of the  $\text{Nb}_3\text{S}_4$  crystals were large enough to allow an investigation of the electrical properties by a four-point-contact method. Metallic behaviour was found with  $\rho \approx 10^{-4} \text{ ohm.cm}$  at room temperature and  $\frac{1}{\rho} \times \frac{d\rho}{dT} = 0.004 \text{ deg}^{-1}$ ; the Seebeck coefficient is of negative sign. From  $100\text{--}370^\circ\text{K}$   $\text{Nb}_3\text{S}_4$  shows nearly temperature-independent paramagnetism with  $\chi_{\text{mol}} = +147 \times 10^{-6} \text{ c.g.s. units per mole of Nb}_3\text{S}_4$  (not corrected for diamagnetism).

### Structure determination

Accurate unit cell dimensions of  $\text{Nb}_3\text{S}_4$  were determined from powder diffractograms taken with  $\text{Cu } K\alpha$  radiation ( $20^\circ\text{C}$ ); silicon powder was used as an internal standard [ $\lambda(\text{Cu } K\alpha_1) = 1.54051 \text{ \AA}$ ;  $\lambda(\text{Cu } K\alpha_2) = 1.54433 \text{ \AA}$ ;  $a(\text{Si}) = 5.4306 \text{ \AA}$ ]. Calculations by least squares gave the values:

$$\begin{aligned} a &= 9.5806 \text{ \AA} \text{ (e.s.d. } 0.0004 \text{ \AA)}; \\ c &= 3.3747 \text{ \AA} \text{ (e.s.d. } 0.0003 \text{ \AA)}; \\ c/a &= 0.35224. \end{aligned}$$

With a unit-cell content of  $\text{Nb}_6\text{S}_8$  the density is calculated as  $5.038 \text{ g.cm}^{-3}$ , which is a reasonable value.

The intensities of the X-ray reflexions were measured with a Nonius automatic single-crystal diffractometer by the  $\theta\text{--}2\theta$  scan method; Zr-filtered  $\text{Mo } K$  radiation was used. The dimensions of the crystal were  $0.43 \times 0.16 \times 0.16 \text{ mm}^3$ . Reliable intensities of 1384 independent reflexions with  $\sin \theta/\lambda < 1.20 \text{ \AA}^{-1}$  were measured; the intensities were corrected for Lorentz and polarization factors and for absorption [ $\mu(\text{Mo}) = 74 \text{ cm}^{-1}$ ].

The Laue group is  $6/m$  and reflexions  $00l$  are systematically absent for  $l = 2n + 1$ ; the space group, therefore, is  $P6_3$  or  $P6_3/m$ . The intensity sequence of reflexions  $hkl'$  was found to be the same as that of  $hk(l' \pm 2)$ , showing that all the atoms lie in planes with  $z = \pm \frac{1}{4}$ . Therefore, the centrosymmetric space group  $P6_3/m$  (no. 176) was adopted; the Nb atoms and the six S(2) atoms were placed in positions  $6(h)$ :  $\pm(x, y, \frac{1}{4}; \bar{y}, x - y, \frac{1}{4}; y - x, \bar{x}, \frac{1}{4})$ , the two S(1) atoms in position  $2(c)$ :  $\pm(\frac{1}{3}, \frac{2}{3}, \frac{1}{4})$ . The same arrangement had been chosen for  $\text{Nb}_3\text{Te}_4$  (Selte & Kjekshus, 1964).

Starting from the atomic parameters of  $\text{Nb}_3\text{Te}_4$ , the structure of  $\text{Nb}_3\text{S}_4$  was refined by least squares (with a program based on the block diagonal approximation). The atomic scattering factors were taken from Moore (1963) and corrected for anomalous dispersion (International Tables, 1962).

Extinction effects appeared to be appreciable, so to account for them the seven strongest reflexions (which also showed the greatest difference between  $F_o$  and  $F_c$ ) were excluded in the last cycles of the refinement. The remaining 1377 reflexions were divided into twenty groups of increasing  $F_o$ . To the reflexions of each group a weight was assigned which was inversely proportional to  $\langle (\Delta F)^2 \rangle$ , where the average was taken over all the reflexions of each group. At the end of the re-

Table 1. Observed and calculated structure factors of Nb<sub>3</sub>S<sub>4</sub> on 10 times absolute scale  
 Reflexions marked by an asterisk were not included in the refinement.

H K L				FO				FC				AC				BC																
1	0	0	0	449	410	-810	-15	*	11	8	0	0	143	147	147	10	5	3	1	1204	1295	1294	-43	1	15	1	257	257	-256	-20		
1	0	0	0	301	475	474	1		12	8	0	0	242	241	-240	-15		6	3	1	657	725	-724	-30	2	15	1	85	77	77	5	
1	0	0	0	251	248	248	1		13	8	0	0	181	184	-184	-15		7	3	1	141	139	-139	-2	3	15	1	187	185	-185	-14	
1	0	0	0	449	475	-474	-1		0	9	0	0	302	312	-312	-15		8	1	1	202	201	-200	-1	4	15	1	61	58	-58	-4	
1	0	0	0	677	600	-600	-24		1	9	0	0	177	170	-179	-9		9	3	1	105	90	-90	-5	6	15	1	87	94	94	10	
1	0	0	0	316	438	438	13		2	9	0	0	694	792	791	35		10	3	1	268	272	-272	-10	7	15	1	149	139	138	8	
1	0	0	0	597	652	-652	-12		3	9	0	0	220	225	-218	-11		11	3	1	367	367	366	20	8	15	1	134	128	128	9	
1	0	0	0	620	700	707	30		4	9	0	0	901	939	931	26		12	3	1	43	66	-66	-4	9	15	1	109	111	110	12	
1	0	0	0	88	84	-84	-3		5	9	0	0	220	225	-218	-11		13	3	1	367	367	366	20	10	15	1	121	119	-118	-10	
1	0	0	0	134	133	132	8		6	9	0	0	98	97	-97	-5		14	3	1	229	222	-221	-10	11	15	1	239	239	230	-21	
1	0	0	0	252	245	245	3		7	9	0	0	288	285	-284	-10		15	3	1	51	40	40	3	12	15	1	352	360	-358	-31	
1	0	0	0	310	315	-315	-14		8	9	0	0	223	235	-214	-12		16	3	1	192	184	-184	-14	13	15	1	251	253	-252	-19	
1	0	0	0	136	134	-127	-11		9	9	0	0	182	170	-179	-14		17	3	1	117	125	-125	-12	14	15	1	100	101	101	8	
1	0	0	0	170	161	160	13		10	9	0	0	244	243	242	19		18	3	1	352	419	-419	-12	15	15	1	68	58	-58	-9	
1	0	0	0	161	157	156	11		11	9	0	0	182	181	-180	-13		19	3	1	375	389	389	11	16	15	1	186	185	185	14	
1	0	0	0	178	170	-170	-12		12	9	0	0	319	333	332	27		20	3	1	764	928	-928	-31	17	15	1	121	119	-118	-10	
1	0	0	0	123	134	133	13		13	9	0	0	92	86	86	7		21	3	1	131	126	-126	-10	18	15	1	352	360	-358	-31	
1	0	0	0	408	410	-410	-15	*	14	9	0	0	359	333	132	6		22	3	1	275	269	-269	-8	19	15	1	485	522	-522	-19	
1	0	0	0	597	652	-652	-12		15	9	0	0	375	385	-385	-18		23	3	1	484	73	73	3	20	15	1	327	361	-361	-14	
1	0	0	0	516	692	692	18		16	9	0	0	192	195	-194	-9		24	3	1	268	269	269	14	21	15	1	172	157	-157	-1	
1	0	0	0	920	920	-920	-20		17	9	0	0	75	74	46	3		25	3	1	304	304	304	6	22	15	1	215	215	-215	-17	
1	0	0	0	926	1533	1533	43	*	18	9	0	0	168	171	-171	-9		26	3	1	402	416	416	21	23	15	1	659	732	732	24	
1	0	0	0	592	703	-703	-21		19	9	0	0	594	631	630	38		27	3	1	593	621	-620	-34	24	15	1	371	372	-372	-12	
1	0	0	0	140	138	-138	-2		20	9	0	0	140	138	-138	-2		28	3	1	44	44	-44	-4	25	15	1	378	396	376	22	
1	0	0	0	163	163	-162	-8		21	9	0	0	261	258	257	19		29	3	1	113	105	-105	-9	26	15	1	390	617	-610	-29	
1	0	0	0	252	221	-221	-10		22	9	0	0	54	45	45	4		30	3	1	77	71	71	3	27	15	1	93	74	74	9	
1	0	0	0	920	920	-920	-20		23	9	0	0	57	50	-50	-5		31	3	1	319	321	320	29	28	15	1	130	118	-117	-8	
1	0	0	0	575	626	625	26		24	9	0	0	177	177	-177	-14		32	3	1	193	206	205	20	29	15	1	39	40	-40	-3	
1	0	0	0	102	107	-107	-8		25	9	0	0	95	94	94	7		33	3	1	232	234	234	10	30	15	1	234	234	234	14	
1	0	0	0	530	562	561	33		26	9	0	0	49	45	45	3		34	3	1	100	10	10	-1	31	15	1	292	284	284	16	
1	0	0	0	92	94	-94	-3		27	9	0	0	36	29	29	4		35	3	1	667	678	678	14	32	15	1	317	301	-301	-10	
1	0	0	0	278	250	250	13		28	9	0	0	111	63	68	65	32		36	3	1	402	413	-413	-12	33	15	1	119	118	117	11
1	0	0	0	139	128	-127	-8		29	9	0	0	214	207	-207	-13		37	3	1	279	272	-271	-12	34	15	1	157	149	-148	-13	
1	0	0	0	156	156	-155	-10		30	9	0	0	90	89	89	8		38	3	1	549	572	-572	-13	35	15	1	327	341	-341	-12	
1	0	0	0	171	172	-171	-14		31	9	0	0	51	19	-18	-5		39	3	1	140	139	139	5	36	15	1	156	159	-159	-12	
1	0	0	0	41	21	21	3		32	9	0	0	190	149	-147	-11		40	3	1	171	166	165	6	37	15	1	111	126	126	-13	
1	0	0	0	129	132	-132	-10		33	9	0	0	98	88	88	7		41	3	1	170	170	170	8	38	15	1	318	324	-324	-14	
1	0	0	0	495	474	474	12		34	9	0	0	178	176	175	12		42	3	1	400	410	409	22	39	15	1	181	179	178	5	
1	0	0	0	304	315	-315	-15		35	9	0	0	184	186	-186	-14		43	3	1	114	129	-129	-14	40	15	1	424	424	424	18	
1	0	0	0	411	477	-477	-12		36	9	0	0	131	131	-131	-12		44	3	1	180	175	-174	-13	41	15	1	541	548	548	17	
1	0	0	0	297	314	-313	-8		37	9	0	0	260	261	-261	-14		45	3	1	237	235	235	16	42	15	1	989	1212	-1211	-42	
1	0	0	0	4	40	40	4		38	9	0	0	32	23	23	2		46	3	1	420	427	427	13	43	15	1	244	244	244	10	
1	0	0	0	343	354	354	10		39	9	0	0	221	216	-216	-14		47	3	1	59	54	-54	-2	44	15	1	128	123	123	2	
1	0	0	0	54	606	606	23		40	9	0	0	40	40	40	27		48	3	1	168	168	168	16	45	15	1	198	198	198	10	
1	0	0	0	43	43	-43	-3		41	9	0	0	32	23	23	2		49	3	1	354	362	-362	-12	46	15	1	451	465	464	20	
1	0	0	0	591	893	892	35		42	9	0	0	71	55	54	5		50	3	1	304	308	308	6	47	15	1	951	957	-957	-25	
1	0	0	0	289	288	-288	-11		43	9	0	0	244	249	-249	-18		51	3	1	540	572	572	24	48	15	1	96	79	79	7	
1	0	0	0	164	157	-157	-10		44	9	0	0	173	171	-171	-12		52	3	1	38	28	28	2	49	15	1	500	507	-506	-32	
1	0	0	0	9	294	-294	-15		45	9	0	0	44	44	44	4		53	3	1	612	668	-668	-24	50	15	1	302	314	-314	-14	
1	0	0	0	231	232	-231	-11		46	9	0	0	114	118	-117	-11		54	3	1	447	477	-477	-20	51	15	1	208	199	-199	-12	
1	0	0	0	304	315	-315	-15		47	9	0	0	31	25	-23	-8		55	3	1	254	248	-246	-10	52	15	1	118	116	116	7	
1	0	0	0	197	201	-201	-14		48	9	0	0	83	79	79	5		56	3	1	45	35	35	0	53	15	1	149	144	143	12	
1	0	0	0	311	307	306	26		49	9	0	0	279	276	-276	-15		57	3	1	365	370	-369	-20	54	15	1	136	160	159	14	
1	0	0	0	113	113	-113	-8		50	9	0	0	184	175	-174	-13		58	3	1	44	49	-49	-4	55	15	1	174	21	-21	-17	
1	0	0	0	252	248	248	1		51	9	0	0	178	177	-177	-13		59	3	1	297	297	296	19</								

Table 1 (cont.)

Table with 16 columns: H, K, L, FC, FC, AC, BC, H, K, L, FC, FC, AC, BC, H, K, L, FC, FC, AC, BC. The table contains numerical data organized in a grid-like structure.

Table 1 (cont.)

H	K	L	FO	FC	AC	BC	H	K	L	FO	FC	AC	BC	H	K	L	FO	FC	AC	BC	H	K	L	FO	FC	AC	BC	
5	1	6	251	240	239	16	0	5	6	301	294	294	19	5	9	6	104	98	88	8	3	3	7	204	213	-212	-17	
6	1	6	54	48	48	1	0	5	6	243	237	-236	-14	6	9	6	92	44	44	4	4	3	7	96	101	101	8	
7	1	6	66	64	63	7	2	5	6	72	70	70	8	6	10	6	45	62	-62	-7	5	3	7	362	378	-377	-30	
8	1	6	98	92	91	8	3	5	6	233	226	226	15	1	10	6	81	86	86	14	6	1	7	214	233	232	22	
9	1	6	232	210	210	16	4	5	6	85	87	87	7	2	10	6	167	170	170	14	7	3	7	35	35	35	1	
10	1	6	259	240	239	-10	5	6	295	270	-270	-14	6	10	6	81	86	86	14	6	1	7	97	100	99	8		
12	1	6	259	249	-248	-26	6	5	6	209	207	206	15	2	11	6	297	300	-299	-24	2	4	7	248	258	257	23	
13	1	6	107	107	-107	-10	7	6	307	317	-317	-31	3	11	6	104	103	103	10	7	4	7	67	67	67	7		
1	2	6	212	202	202	15	9	5	6	107	111	-110	-11	2	11	6	125	119	119	11	6	4	7	82	85	-85	-9	
2	2	6	198	132	132	9	0	6	6	166	158	-155	-10	0	12	6	105	110	109	11	7	4	7	97	102	-101	-11	
3	2	6	113	98	98	6	1	6	6	261	260	-259	-17	1	13	6	145	144	143	12	8	5	7	66	69	-69	-9	
4	2	6	58	15	-15	2	2	6	6	79	62	62	4	0	13	6	145	144	143	12	9	5	7	66	69	-69	-9	
5	2	6	137	227	-227	-18	3	6	6	147	144	-143	-24	1	14	6	151	152	-151	-13	2	5	7	180	198	-198	-18	
6	2	6	353	350	-358	-27	4	6	6	168	172	171	14	2	15	6	118	118	-115	-17	3	5	7	158	158	157	14	
7	2	6	242	245	244	19	5	6	6	171	172	171	14	3	16	6	144	155	-154	-13	4	5	7	86	87	86	8	
8	2	6	79	75	75	8	6	6	6	118	116	115	9	4	17	6	100	-99	-9	9	6	5	7	46	4	-4	0	
9	2	6	129	133	132	11	7	6	6	142	151	-150	-14	5	18	6	70	60	60	9	7	6	7	178	193	192	18	
10	3	6	75	63	-63	0	8	7	6	255	255	254	15	6	19	6	272	278	277	24	8	6	7	49	49	48	5	
1	3	6	407	404	-403	-26	9	7	6	148	144	143	11	7	20	6	113	105	-105	-7	9	6	7	105	105	105	6	
2	3	6	70	24	0	0	0	7	6	67	52	52	3	8	21	6	92	55	54	9	1	7	7	79	86	-86	-6	
3	3	6	66	50	-50	-3	1	7	6	214	206	205	15	0	1	7	149	152	151	13	2	8	7	177	194	-193	-19	
4	3	6	285	279	278	19	2	7	6	158	152	-152	-13	1	2	7	65	55	54	9	3	7	7	265	278	-277	-24	
5	3	6	177	175	174	14	3	7	6	146	141	-140	-12	2	1	7	92	68	68	5	4	6	7	177	194	-193	-19	
6	3	6	122	119	119	8	4	7	6	102	102	-102	-16	3	1	7	81	88	88	6	5	7	7	137	140	-140	-12	
7	3	6	152	119	-117	-10	5	7	6	152	120	-120	-12	4	1	7	200	205	-204	-18	6	7	7	198	208	207	19	
8	3	6	120	134	-133	-14	6	8	6	284	278	-277	-23	6	1	7	85	87	86	8	7	7	7	119	123	122	10	
9	3	6	88	93	92	8	7	6	6	46	42	42	3	7	1	7	104	104	-104	-10	8	7	7	105	108	107	19	
10	4	6	214	202	202	13	8	6	6	156	156	-156	-9	8	1	7	104	207	206	19	9	7	7	118	125	125	10	
1	4	6	52	19	18	1	9	6	6	91	88	88	6	9	2	7	118	116	115	7	2	8	7	100	104	-104	-10	
2	4	6	312	301	300	20	0	7	6	139	145	145	12	1	2	7	362	389	-387	-30	3	7	7	265	278	-277	-24	
3	4	6	289	274	-275	-22	1	8	6	6	77	86	-87	-8	2	2	7	302	308	307	24	4	7	7	213	221	220	20
4	4	6	93	94	92	7	2	8	6	127	132	131	11	3	2	7	151	155	155	14	5	8	7	409	442	441	36	
5	4	6	274	269	-268	-21	3	0	7	0	43	31	2	4	2	7	117	119	118	10	6	0	8	139	109	-109	-10	
6	4	6	92	92	92	7	1	9	6	82	79	79	7	5	2	7	49	43	43	3	7	0	8	96	76	76	8	
7	4	6	27	69	69	6	2	9	6	321	331	-330	-27	6	2	7	101	101	-101	10	8	1	8	101	109	-109	-10	
8	4	6	93	100	100	11	3	9	6	131	131	131	11	7	0	7	150	155	154	13	9	1	8	86	33	-33	-13	
9	4	6	137	140	139	11	4	9	6	258	251	-250	-22	8	1	7	195	197	-196	-14	10	2	8	86	76	76	8	

finement the discrepancy factor  $R = \frac{\sum |F_o - F_c|}{\sum |F_o|}$  was 0.06. Observed and calculated structure factors are listed in Table 1.

The final positional and thermal parameters with the standard deviations calculated by the least-squares program are given in Table 2. The structure was checked by a difference Fourier synthesis, which showed no unusual features. In particular, no evidence of scattering was found in the large channels parallel to the *c* axis around the origin (Fig. 1); this shows that the adopted composition Nb<sub>3</sub>S<sub>4</sub> is correct.

Table 2. Positional and thermal parameters of Nb<sub>3</sub>S<sub>4</sub> (Å<sup>2</sup>)

Standard deviations in units of the last decimal are given in brackets.

	Nb	S(1)	S(2)
<i>x</i>	0.48538 (2)	$\frac{1}{3}$	0.34495 (7)
<i>y</i>	0.11392 (2)	$\frac{2}{3}$	0.28723 (7)
<i>z</i>	$\frac{1}{4}$	$\frac{1}{4}$	$\frac{1}{4}$
<i>U</i> <sub>11</sub>	0.00400 (5)	0.00449 (12)	0.00711 (15)
<i>U</i> <sub>22</sub>	0.00340 (5)	0.00449 (12)	0.00544 (15)
<i>U</i> <sub>33</sub>	0.00399 (5)	0.00418 (24)	0.00500 (16)
2 <i>U</i> <sub>12</sub>	0.00355 (8)	0.00449 (12)	0.00719 (26)
2 <i>U</i> <sub>13</sub> = 2 <i>U</i> <sub>23</sub>	0	0	0

### Discussion

The structure of Nb<sub>3</sub>S<sub>4</sub> is shown in Fig. 1; interatomic distances and their standard deviations are listed in Table 3.

The structure is built up by distorted NbS<sub>6</sub> octahedra joined together to form a three-dimensional network; each octahedron is linked to two other octahedra by common faces and to four octahedra by common edges. The Nb atoms are not at the centres of the octahedra but are shifted by 0.306 Å in the direction of a face sharing two edges with other octahedra; thus zigzag Nb–Nb–Nb chains are formed running in the *c* direction (which is the needle axis). The metal–metal

Table 3. Interatomic distances in Nb<sub>3</sub>S<sub>4</sub> (Å)

Standard deviations in units of the last decimal are given in brackets.

Nb–2 S(2)	2.454 (1)	Nb–2 Nb	2.8813 (6)
–1 S(2)	2.492 (1)	–2 Nb	3.3695 (6)
–2 S(1)	2.5752 (3)	–2 Nb	3.3747 (3)
–1 S(2)	2.608 (1)		
S(1)–6 Nb	2.5752 (3)	S(1)–6 S(2)	3.344 (1)
		–2 S(1)	3.3747 (3)
		–3 S(2)	3.688 (1)
S(2)–2 Nb	2.454 (1)	S(2)–2 S(1)	3.344 (1)
–1 Nb	2.492 (1)	–2 S(2)	3.3747 (3)
–1 Nb	2.608 (1)	–4 S(2)	3.498 (2)
		–1 S(1)	3.688 (1)
		–2 S(2)	4.020 (2)

distances (and also the next-nearest-neighbour distances) in the chains are comparable to those in metallic Nb (*cf.* Table 4). Similar zigzag metal–metal chains have been found in Mo<sub>2</sub>S<sub>3</sub> (Jellinek, 1961) and the isotypic compounds Nb<sub>2</sub>Se<sub>3</sub> and Ta<sub>2</sub>Se<sub>3</sub> (Kadijk, Huisman & Jellinek, 1968), in Nb<sub>3</sub>Se<sub>4</sub> and Nb<sub>3</sub>Te<sub>4</sub> (Selte & Kjekshus, 1964) and in the MnP-type substructure of TiSe<sub>1-x</sub> (Grønvold & Langmyhr, 1961) and probably also of orthorhombic NbS<sub>1-x</sub> (Kadijk & Jellinek, 1968). Some Nb–Nb distances are compared in Table 4.

Another remarkable feature of the Nb<sub>3</sub>S<sub>4</sub> structure is the number of large infinite channels in the *c* direction (Fig. 1). Although these channels are wide enough to accommodate additional atoms [distance S(2) to *c* axis 2.807 Å; Nb to *c* axis 3.722 Å], they appear to be empty. Thus, S(2) has very asymmetric surroundings formed by four Nb atoms, while S(1) has six Nb neighbours in a trigonal prism. The relation between this structure and those of 2*s*–Nb<sub>1+x</sub>S<sub>2</sub> and Nb<sub>5</sub>Se<sub>4</sub> has been discussed by Selte & Kjekshus (1964).

The computer programs used in this investigation were written by Mrs R. Olthof-Hazekamp, Dr D. W. Smits, Dr J. H. Palm and Ir W. J. A. M. Peterse, and Mr. J. C. Eikelenboom.

Table 4. *Metal-metal distances (Å) in some Nb chalcogenides and in Nb metal*

	Nb <sub>3</sub> S <sub>4</sub>	Nb <sub>2</sub> Se <sub>3</sub> *		Nb <sub>3</sub> Te <sub>4</sub> †	Nb
		Nb(1)	Nb(2)		
Within zigzag chains					
Nb-Nb	2.8813	2.972	3.127	2.973	2.8585 (8 ×)
Nb-Nb	3.3747	3.434	3.434	3.647	3.3007 (6 ×)
∠ Nb-Nb-Nb	71.69°	70.57°	66.62°	75.66°	70.53°
Between chains					
Nb-Nb	3.3695 (2 ×)	3.329 (1 ×)	3.329 (1 ×)	3.854 (2 ×)	

\* Kadijk, Huisman &amp; Jellinek (1968).

† Selte &amp; Kjekshus (1964).

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## Zur Struktur des Wasserstoffperoxids

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Conflicting results of space group determinations of hydrogen peroxide in the literature have been checked by X-ray photographs of well formed hydrogen peroxide single crystals, using the Weissenberg camera. The observed reflexions 100, 001, 002, 003, 006 and 007, which are not allowed by the space group  $P4_12_1$ , originate from double reflexions of the pairs of planes  $(\bar{2}21)-(3\bar{2}\bar{1})$ ,  $(11\bar{3})-(\bar{1}\bar{1}4)$ ,  $(120)-(\bar{1}\bar{2}2)$ ,  $(123)-(\bar{1}\bar{2}0)$ ,  $(\bar{1}22)-(\bar{1}24)$ , and  $(336)-(\bar{3}\bar{3}\bar{1})$ . The space group  $P4_12_1$  is confirmed. This is a further interesting example of 'the double reflexion' effect as a possible source of error in structure determination.

Die Ergebnisse einer vollständigen röntgenographischen Feinstrukturuntersuchung an Wasserstoffperoxid wurden von Abrahams, Collin & Lipscomb (1951) mitgeteilt. Danach kristallisiert die Substanz im tetragonalen Kristallsystem mit den Gitterkonstanten  $a = 4,06 \pm 0,02$  und  $c = 8,00 \pm 0,02$  Å. Als Raumgruppe wird  $P4_12_1$  genannt. Diese Angaben stehen teilweise im Widerspruch zu früheren Arbeiten von Fehér & Klötzer (1935, 1937), die zwar ebenfalls die Gitterkonstanten zu  $a = 4,02$  und  $c = 8,02$  Å bestimmten, aber auf Grund fehlender systematischer Auslöschungen und der Laue-Symmetrie  $D_4^1$ ,  $D_{4h}^1$ ,  $C_4^1$ ,  $D_{2d}^2$  und  $D_{2d}^3$  als wahrscheinliche Raumgruppen anführten. Da die Ergebnisse von Fehér & Klötzer aus Untersuchungen an

freien Wasserstoffperoxid-Einkristallen gewonnen wurden, während Abrahams *et al.* ihre Arbeiten an Wasserstoffperoxid ausgeführt hatten, das in Pyrexglaskapillaren kristallin erstarrt war, bestand die Möglichkeit, dass die abweichenden Angaben durch diese voneinander verschiedenen Arten der recht schwierigen Herstellung und Präparation der Einkristalle bedingt war.

Um die unterschiedlichen Angaben hinsichtlich der Raumgruppe zu überprüfen, haben wir die röntgenographischen Untersuchungen an freien Wasserstoffperoxid-Einkristallen erneut aufgegriffen.

Gut ausgebildete nadelförmige Einkristalle von Wasserstoffperoxid erhält man durch Animpfen von